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## Review

## Typical Notation



## Positivity, Self-Adjointness

$$
\left\{f(\cdot) \in C^{\infty}(\Omega):\left.f\right|_{\partial \Omega} \equiv 0\right\}
$$



$$
\begin{aligned}
\mathcal{L}[f] & :=-\Delta f \\
\langle f, g\rangle & :=\int_{\Omega} f(x) g(x) d x
\end{aligned}
$$

On board:
I. Positive: $\langle f, \mathcal{L}[f]\rangle \geq 0$
2. Self-adjoint: $\langle f, \mathcal{L}[g]\rangle=\langle\mathcal{L}[f], g\rangle$

## Proof

## Proof of 1

$\langle f, \mathscr{L}[f]\rangle=\int_{\Omega} f(-\nabla \cdot \nabla f) d V=\int_{\partial \Omega} f(-\nabla f \cdot \vec{n}) d S+\int_{\Omega} \nabla f \cdot \nabla f d V=\int_{\Omega} \nabla f \cdot \nabla f d V \geq 0$
where the second equality follows from Green formula, and the third equality follows from $\left.f\right|_{\partial \Omega} \equiv 0$

## Proof of 2

$\langle f, \mathscr{L}[g]\rangle=\int_{\Omega} f(-\nabla \cdot \nabla g) d V=\int_{\partial \Omega} f(-\nabla g \cdot \vec{n}) d S+\int_{\Omega} \nabla f \cdot \nabla g d V=\int_{\Omega} \nabla f \cdot \nabla g d V$
where the second equality follows from Green formula, and the third equality follows from $\left.f\right|_{\partial \Omega} \equiv 0$
Similarly, $\langle\mathscr{L}[f], g\rangle=\int_{\Omega} \nabla g \cdot \nabla f d V$

It also shows $\langle f, \mathscr{L}[g]\rangle=\int_{\Omega} \nabla f \cdot \nabla g d V$

## Laplacian(-Bertrami) Operator Diagonalizable!

## Theorem. Let $\boldsymbol{H} \neq \mathbf{0}$ be an infinite-dimensional,

 separable Hilbert space and let $K \in L(H)$ be compact and self-adjoint. Then, there exists a countable orthonormal basis of $\boldsymbol{H}$ consisting of eigenvectors of $\boldsymbol{K}$.

Hilbert space: Space with inner product
Separable: Admits countable, dense subset
Compact operator: Bounded sets to relatively compact sets
Self-adjoint: $\langle K v, w\rangle=\langle v, K w\rangle$

## Eigenhomers



## Dirichlet Energy

$$
E[f]:=\int_{\Omega}\langle\nabla f, \nabla f\rangle d A
$$

## Proof

We use variational method to derive.
Lagrangian: $\quad \mathbb{L}[f]=\frac{1}{2} \int\langle\nabla f, \nabla f\rangle+\int_{\partial \Omega} \lambda(x)(f(x)-g(x))$
So
$\delta \mathbb{L}[f]=\mathbb{C}[f+\delta h]-\mathbb{L}[f]=\int_{\Omega}\langle\nabla f, \nabla \delta h\rangle+\int_{\partial \Omega} \lambda(x) \delta h(x)=\left\{\int_{\partial \Omega} \delta h(\nabla f \cdot \vec{n})-\int_{\Omega} \delta h(\nabla \cdot \nabla f)\right\}+\int_{\partial \Omega} \lambda(x) \delta h(x)$
In the interior of $\Omega, \Delta f \equiv 0$ so that $\delta \mathbb{\square}[f]=0$ for any $\delta h$
Note: in the derivation we ignored the second-order infinitesimal term $\mathrm{O}\left(\|\delta h\|^{2}\right)$

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# CSE291-C00 Discrete Laplacian \& Its Applications 

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## Discrete Laplacian

## Our Focus

## $f \in C^{\infty}(M) \longrightarrow \leadsto \Delta f \in C^{\infty}(M)$

## The Laplacian

## Recall: Planar Region



Wave equation:

$$
\frac{\partial^{2} u}{\partial t^{2}}=\Delta u
$$

$$
\Delta:=\sum_{i} \frac{\partial^{2}}{\partial x_{i}^{2}}
$$

## Discretizing the Laplacian



## Problem

## Laplacian is a differential operator!



## A Principled Approach to Connect Continuous \& Discrete Objects

First-order Galerkin

# Finite element method (FEM) 



## Integration by Parts to the Rescue

$\int_{\Omega} f \Delta g d A=$ boundary terms $-\int_{\Omega} \nabla f \cdot \nabla g d A$

a GUIDE to
INTEGRATION BY PARTS:
GIVEN A PROBLEM OF THE FORM:

$$
\int f(x) g(x) d x=?
$$

CHOOSE VARIABLES $U$ AND $V$ SUCH THAT:

$$
\begin{aligned}
u & =f(x) \\
d v & =g(x) d x
\end{aligned}
$$

NOW THE ORIGINAL EXPRESSION BECONES:

$$
\int u d v=?
$$

WHICH DEFINITELY LOOKS EASIER. ANYWAY, I GOTTA RUN. BUT GOOD LWCK!

## Slightly Easier?

$\int_{\Omega} f \Delta g d A=$ boundary terms $-\int_{\Omega} \nabla f \cdot \nabla g d A$
Laplacian
(second derivative)

## Slightly Easier?

$\int_{\Omega} f \Delta g d A=$ boundary terms $-\int_{\Omega} \nabla f \cdot \nabla g d A$

Kinda-sorta cancels out?

## Overview:Galerkin FEM Approach

$$
\begin{aligned}
& g=\Delta f \\
\Longrightarrow & \int \psi g d A=\int \psi \Delta f d A=-\int(\nabla \psi \cdot \nabla f) d A
\end{aligned}
$$

## L2 Dual of a Function

Function $f: M \rightarrow \mathbb{R}$
$\downarrow$
$\mathcal{L}_{f}: L^{2}(M) \rightarrow \mathbb{R}$
$\mathcal{L}_{f}[g]:=\int_{M} f(x) g(x) d A$
"Test function"

## Observation



Can recover function from dual

## Dual of Laplacian

$$
\begin{aligned}
\{g & \left.\in L^{\substack{\text { Space of test functions: }}}(M):\left.g\right|_{\partial M} \equiv 0\right\} \\
\mathcal{L}_{\Delta f}[g] & =\int_{M} g \Delta f d A \\
& =-\int_{M} \nabla g \cdot \nabla f d A
\end{aligned}
$$

## Use Laplacian without evaluating it!

## Galerkin's Approach

Choose one of each:
-Function space
-Test functions
Often the same!

## One Derivative is Enough

$$
\mathcal{L}_{\Delta f}[g]=-\int_{M} \nabla g \cdot \nabla f d A
$$

## Representing Functions



## What Do We Need

## What Do We Need



## What Do We Need

$$
\mathcal{L}_{\Delta f}[g]=-\int_{M} \nabla g, \nabla f d A
$$

## What Do We Need



Sum scalars per face multiplied by face areas

## Gradient of a Hat Function



## Gradient of a Hat Function



## Gradient of a Hat Function

$$
\underbrace{f\left(v_{1}\right)=1}_{f\left(v_{2}\right)=0}
$$

## Gradient of a Hat Function



## What We Actually Need

$$
\mathcal{L}_{\Delta f}[g]=-\int_{M} \nabla g \cdot \nabla f d A
$$



## What We Actually Need

$$
\mathcal{L}_{\Delta f}[g]=-\int_{M} \nabla g \cdot \nabla f d A
$$

Case I: Same vertex

$$
\begin{aligned}
\int_{T}\langle\nabla f, \nabla f\rangle d A & =A\|\nabla f\| 2 \\
& =\frac{A}{h^{2}}=\frac{b}{2 h} \\
& =\frac{1}{2}(\cot \alpha+\cot \beta)
\end{aligned}
$$

## What We Actually Need

## $\mathcal{L}_{\Delta f}[g]=-\int_{M} \nabla g \cdot \nabla f d A$

## Case 2: Different vertices

$$
\begin{aligned}
\int_{T}\left\langle\nabla f_{\alpha}, \nabla f_{\beta}\right\rangle d A & =A\left\langle\nabla f_{\alpha}, \nabla f_{\beta}\right\rangle \\
& =\frac{1}{4 A}\left\langle e_{31}^{\perp}, e_{12}^{\perp}\right\rangle=-\frac{\ell_{1} \ell_{2} \cos \theta}{4 A} \\
& =\frac{-h^{2} \cos \theta}{4 A \sin \alpha \sin \beta}=\frac{-h \cos \theta}{2 b \sin \alpha \sin \beta} \\
& =-\frac{\cos \theta}{2 \sin (\alpha+\beta)}=-\frac{1}{2} \cot \theta
\end{aligned}
$$

## Summing Around a Vertex



$$
\left\langle\nabla h_{p}, \nabla h_{p}\right\rangle=\frac{1}{2} \sum_{i}\left(\cot \alpha_{i}+\cot \beta_{i}\right)
$$



$$
\left\langle\nabla h_{p}, \nabla h_{q}\right\rangle=\frac{1}{2}\left(\cot \theta_{1}+\cot \theta_{2}\right)
$$

## THE COTANGENT LAPLACIAN

$L_{i j}= \begin{cases}\frac{1}{2} \sum_{i \sim k}\left(\cot \alpha_{i k}+\cot \beta_{i k}\right) & \text { if } i=j \\ -\frac{1}{2}\left(\cot \alpha_{i j}+\cot \beta_{i j}\right) & \text { if } i \sim j \\ 0 & \text { otherwise }\end{cases}$


## Poisson Equation

## $\Delta f=g$



## Eigenhomers



## Point Cloud Laplace: Easiest Option

$$
\begin{aligned}
W_{i j} & =\exp \left(-\frac{\left\|x_{i}-x_{j}\right\|^{2}}{t}\right) \\
D_{i i} & =\sum_{j} W_{j i} \\
L & =D-W \\
L f & =\lambda D f
\end{aligned}
$$

"Laplacian Eigenmaps for Dimensionality Reduction and Data Representation"

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Applications of Laplacian:
Intrinsic Shape Descriptor

## Why Study the Laplacian?

- Encodes intrinsic geometry

Edge lengths on triangle mesh, Riemannian metric on manifold

- Multi-scale

Filter based on frequency

- Geometry through linear algebra Linear/eigenvalue problems, sparse positive definite matrices
- Connection to physics

Heat equation, wave equation, vibration, ...

## Example Task: Shape Descriptors



Pointwise quantity

## Isometry Invariance: Hope



## Descriptor Tasks

- Characterize local geometry Feature/anomaly detection
- Describe point's role on surface Symmetry detection, correspondence


## Descriptors We've Seen Before



## Gaussian and mean curvature

## Desirable Properties

- Distinguishing

Provides useful information about a point

- Stable

Numerically and geometrically

- Intrinsic

No dependence on embedding
Sometimes
undesirable!

## Intrinsic Descriptors

## Invariant under

- Rigid motion
- Bending without stretching


## Intrinsic Descriptor



Theorema Egregium ("Totally Awesome Theorem"):
Gaussian curvature is intrinsic.

## Gaussian curvature

## End of the Story?

$$
K=\kappa_{1} \kappa_{2}
$$

## Second derivative quantity

## End of the Story?


http://www.integrityware.com/images/MerceedesGaussianCurvature.jpg
Non-unique

## Desirable Properties

# Incorporates neighborhood information in an intrinsic fashion 

Stable under small deformation

## Recall: Connection to Physics



## Intrinsic Observation

## Heat diffusion patterns are not affected if you bend a surface.

## Global Point Signature


"Laplace-Beltrami Eigenfunctions for Deformation Invariant Shape Representation" Rustamov, SGP 2007

## Global Point Signature

$$
\begin{gathered}
\text { l } \\
(p):=\left(-\frac{1}{\sqrt{\lambda_{1}}} \phi_{1}(p),-\frac{1}{\sqrt{\lambda_{2}}} \phi_{2}(p),-\frac{1}{\sqrt{\lambda_{3}}} \phi_{3}(p), \cdots\right)
\end{gathered}
$$

If surface does not self-intersect, neither does the GPS embedding.

Proof: Laplacian eigenfunctions span ; if GPS(p)=GPS(q), then all functions on would be equal at $p$ and $q$.

## Global Point Signature



## GPS is isometry-invariant.

Proof: Comes from the Laplacian.

## Recall: Connection to Physics



## Heat Kernel Map



## How much heat diffuses from $p$ to $x$ in time $t$ ?

One Point Isometric Matching with the Heat Kernel
Ovsjanikov et al. 2010

## Heat Kernel Map



$$
\operatorname{HKM}_{p}(x, t):=k_{t}(p, x)
$$

## Theorem: Only have to match one point!

One Point Isometric Matching with the Heat Kernel

