

Lecture 11:
High-Dimensional Geometry

Focus of Past High-D Geometry Research

From Low-D to High-D, Change of View

- From calculus view to statistical view
- From infinitesimal analysis to structural/topological analysis
- We will focus on the “odd” behaviors in high-D geometry

Some basic tools

Theorem 2.1 (Markov's inequality) *Let x be a nonnegative random variable. Then for $a > 0$,*

$$\text{Prob}(x \geq a) \leq \frac{E(x)}{a}.$$

Proof on board

Theorem 2.3 (Chebyshev's inequality) *Let x be a random variable with mean m and variance σ^2 . Then*

$$\text{Prob}(|x - m| \geq a\sigma) \leq \frac{1}{a^2}.$$

Theorem 2.4 (Law of large numbers) *Let x_1, x_2, \dots, x_n be n samples of a random variable x . Then*

$$\text{Prob}\left(\left|\frac{x_1 + x_2 + \dots + x_n}{n} - E(x)\right| > \epsilon\right) \leq \frac{\sigma^2}{n\epsilon^2}$$

Volume & Surface Distribution

Volume of Unit Sphere Goes to Zero

- Sphere volume:

$$V(d) = \int_{S^d} \int_{r=0}^1 r^{d-1} dr d\Omega.$$

At radius r , the surface area of the top of the cone is $r^{d-1} d\Omega$ since the surface area is $d - 1$ dimensional and each dimension scales by r

$$V(d) = \int_{S^d} d\Omega \int_{r=0}^1 r^{d-1} dr = \frac{1}{d} \int_{S^d} d\Omega = \frac{A(d)}{d}$$

Volume of Unit Sphere Goes to Zero

- Computation of $A(d)$
 - Trick: compare the integration of e^{-x^2} in Cartesian and Polar systems

Consider a different integral

$$I(d) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \cdots \int_{-\infty}^{\infty} e^{-(x_1^2 + x_2^2 + \cdots + x_d^2)} dx_d \cdots dx_2 dx_1.$$

First, calculate $I(d)$ by integration in Cartesian coordinates.

$$I(d) = \left[\int_{-\infty}^{\infty} e^{-x^2} dx \right]^d = (\sqrt{\pi})^d = \pi^{\frac{d}{2}}.$$

But

$$I(d) = \int_{S^d} d\Omega \int_0^{\infty} e^{-r^2} r^{d-1} dr.$$

Volume of Unit Sphere Goes to Zero

$$I(d) = \left[\int_{-\infty}^{\infty} e^{-x^2} dx \right]^d = (\sqrt{\pi})^d = \pi^{\frac{d}{2}}.$$

$$I(d) = \int_{S^d} d\Omega \int_0^{\infty} e^{-r^2} r^{d-1} dr.$$

$$\int_0^{\infty} e^{-r^2} r^{d-1} dr = \frac{1}{2} \int_0^{\infty} e^{-t} t^{\frac{d}{2}-1} dt = \frac{1}{2} \Gamma\left(\frac{d}{2}\right)$$

and hence, $I(d) = A(d) \frac{1}{2} \Gamma\left(\frac{d}{2}\right)$

Lemma 2.5 *The surface area $A(d)$ and the volume $V(d)$ of a unit-radius sphere in d dimensions are given by*

$$A(d) = \frac{2\pi^{\frac{d}{2}}}{\Gamma\left(\frac{d}{2}\right)} \quad \text{and} \quad V(d) = \frac{2}{d} \frac{\pi^{\frac{d}{2}}}{\Gamma\left(\frac{d}{2}\right)}.$$

Volume of Unit Sphere Goes to Zero

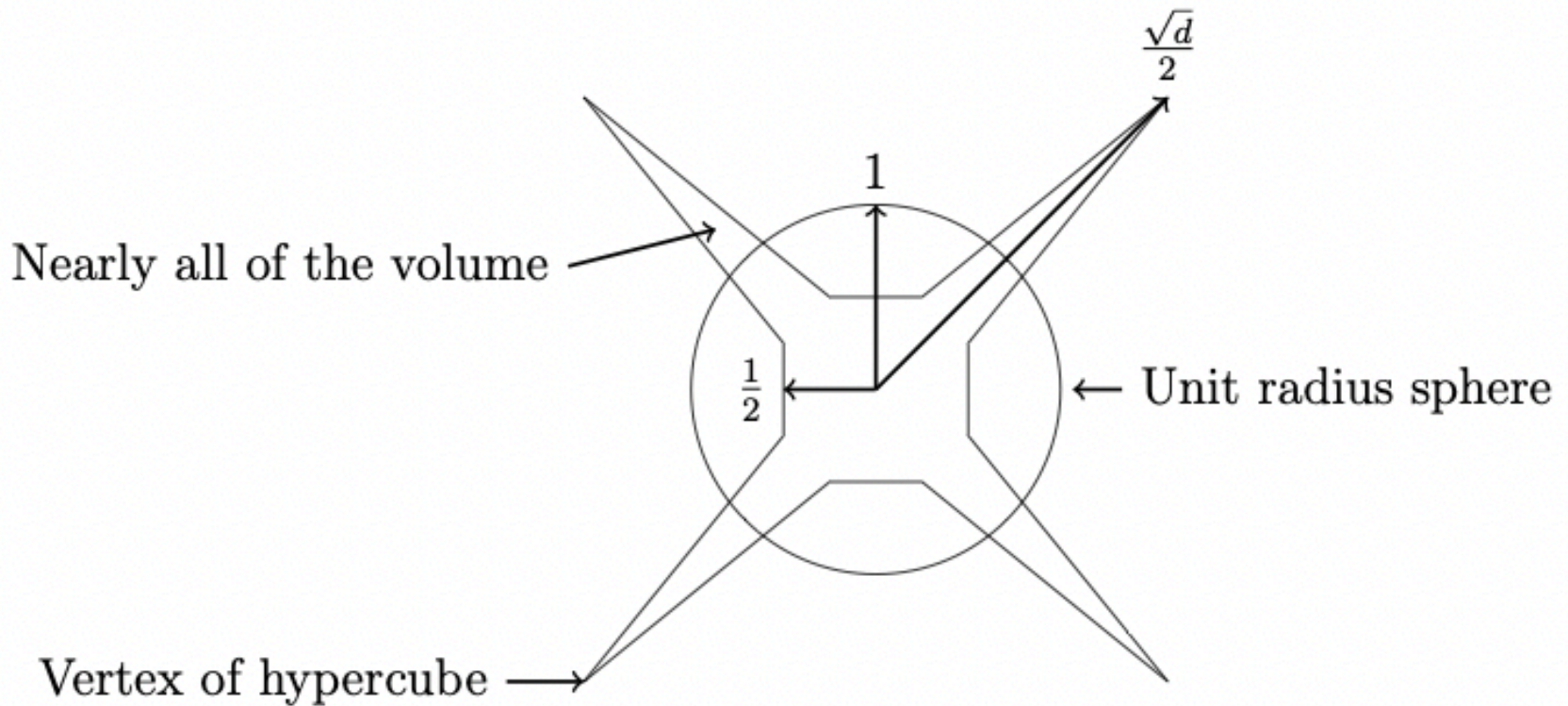
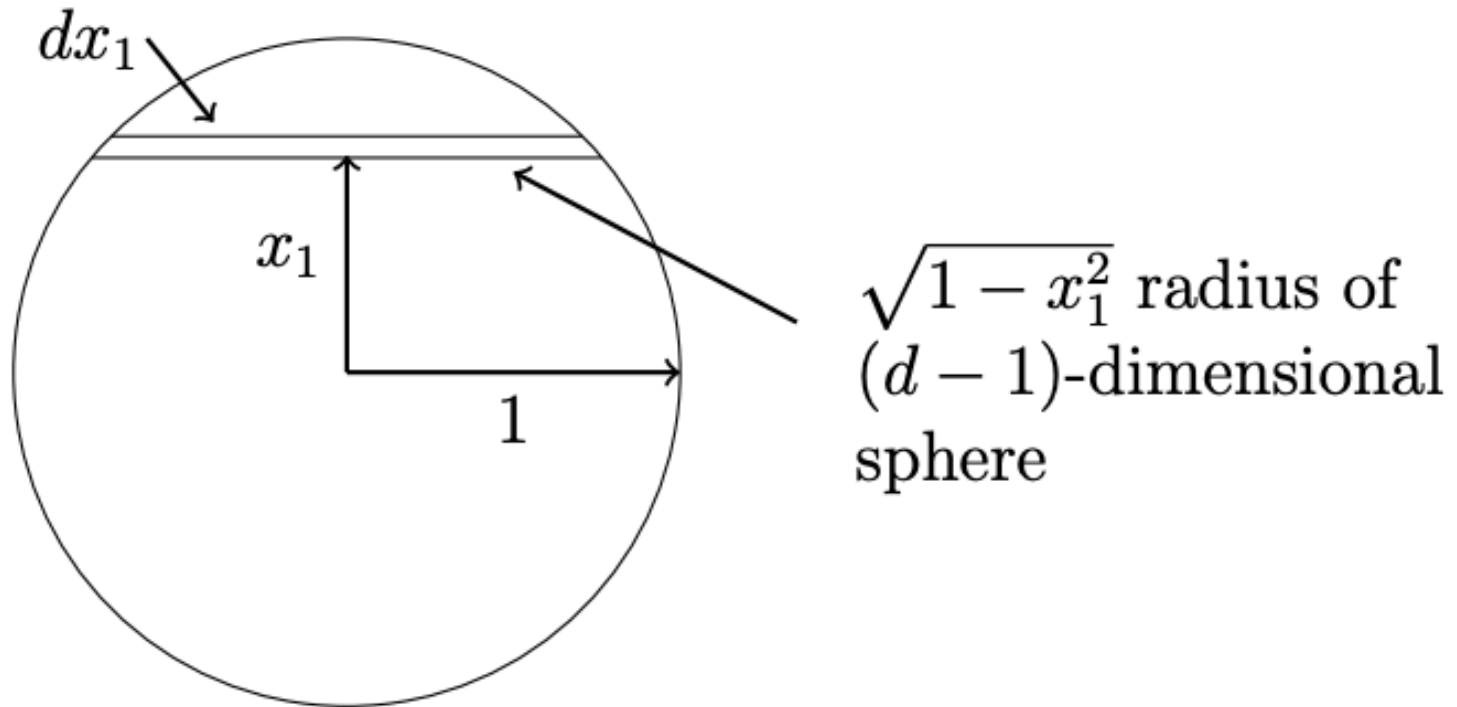


FIGURE 2.5: Conceptual drawing of a sphere and a cube.

The Volume is Near the Equator

- It turns out that essentially all of the volume of the upper hemisphere lies between the plane $x_1 = 0$ and a parallel plane, $x_1 = \varepsilon$, that is slightly higher.
- For what value of ε does essentially all the volume lie between $x_1 = 0$ and $x_1 = \varepsilon$? The answer depends on the dimension. For dimension d , it is $O\left(\frac{1}{\sqrt{d-1}}\right)$

The Volume is Near the Equator



The Volume is Near the Equator

$$\frac{\text{Volume above slice}}{\text{Volume upper hemisphere}} \leq \frac{\text{Upper bound volume above slice}}{\text{Lower bound volume upper hemisphere}}$$

Let $T = \{\mathbf{x} \mid |\mathbf{x}| \leq 1, x_1 \geq \varepsilon\}$ be the portion of the sphere above the slice.

$$\text{Volume}(T) = \int_{\varepsilon}^1 (1 - x_1^2)^{\frac{d-1}{2}} V(d-1) dx_1 = V(d-1) \int_{\varepsilon}^1 (1 - x_1^2)^{\frac{d-1}{2}} dx_1.$$

use the inequality $1 + x \leq e^x$ for all real x

$$\begin{aligned} \text{Volume}(T) &\leq V(d-1) \int_{\varepsilon}^{\infty} e^{-\frac{d-1}{2}x_1^2} dx_1 \\ &\leq V(d-1) \int_{\varepsilon}^{\infty} \frac{x_1}{\varepsilon} e^{-\frac{d-1}{2}x_1^2} dx_1. \end{aligned}$$

$$\text{Volume}(T) \leq \frac{1}{\varepsilon(d-1)} e^{-\frac{d-1}{2}\varepsilon^2} V(d-1)$$

The Volume is Near the Equator

Approximate Volume of upper hemisphere:

volume of the entire upper hemisphere. Clearly, the volume of the upper hemisphere is at least the volume between the slabs $x_1 = 0$ and $x_1 = \frac{1}{\sqrt{d-1}}$, which is at least the volume of the cylinder of radius $\sqrt{1 - \frac{1}{d-1}}$ and height $\frac{1}{\sqrt{d-1}}$. The volume of the cylinder is $1/\sqrt{d-1}$ times the $d-1$ -dimensional volume of the disk $R = \left\{ \mathbf{x} \mid |\mathbf{x}| \leq 1; x_1 = \frac{1}{\sqrt{d-1}} \right\}$. Now R is a $d-1$ -dimensional sphere of radius $\sqrt{1 - \frac{1}{d-1}}$ and so its volume is

$$\text{Volume}(R) = V(d-1) \left(1 - \frac{1}{d-1}\right)^{(d-1)/2}.$$

Using $(1-x)^a \geq 1-ax$

$$\text{Volume}(R) \geq V(d-1) \left(1 - \frac{1}{d-1} \frac{d-1}{2}\right) = \frac{1}{2}V(d-1).$$

Thus, the volume of the upper hemisphere is at least $\frac{1}{2\sqrt{d-1}}V(d-1)$.

The Volume is Near the Equator

Lemma 2.6 *For any $c > 0$, the fraction of the volume of the unit hemisphere above the plane $x_1 = \frac{c}{\sqrt{d-1}}$ is less than $\frac{2}{c}e^{-c^2/2}$.*

Proof: Substitute $\frac{c}{\sqrt{d-1}}$ for ε in the above. ■

The Volume is in a Narrow Annulus

The ratio of the volume of a sphere of radius $1 - \varepsilon$ to the volume of a unit sphere in d -dimensions is

$$\frac{(1 - \varepsilon)^d V(d)}{V(d)} = (1 - \varepsilon)^d,$$

and thus goes to zero as d goes to infinity when ε is a fixed constant. In high dimensions, all of the volume of the sphere is concentrated in a narrow annulus at the surface.

Since, $(1 - \varepsilon)^d \leq e^{-\varepsilon d}$, if $\varepsilon = \frac{c}{d}$, for a large constant c , all but e^{-c} of the volume of the sphere is contained in a thin annulus of width c/d . The important item to remember is that most of the volume of the d -dimensional unit sphere is contained in an annulus of width $O(1/d)$ near the boundary. If the sphere is of radius r , then for sufficiently large d , the volume is contained in an annulus of width $O\left(\frac{r}{d}\right)$.

Gaussian in High Dimension

Gaussian Distribution

$$p(\mathbf{x}) = \frac{1}{(2\pi)^{d/2} \sigma^d} \exp\left(-\frac{|\mathbf{x}|^2}{2\sigma^2}\right)$$

Expected squared distance of a point from the center of a Gaussian

- A 1-dimensional Gaussian has its mass close to the origin.

Expected squared distance of a point from the center of a Gaussian

- However, as the dimension is increased something different happens
- When $\sigma^2 = 1$, integrating the probability density over a unit sphere centered at the origin yields nearly zero mass since the volume of a unit sphere is negligible.

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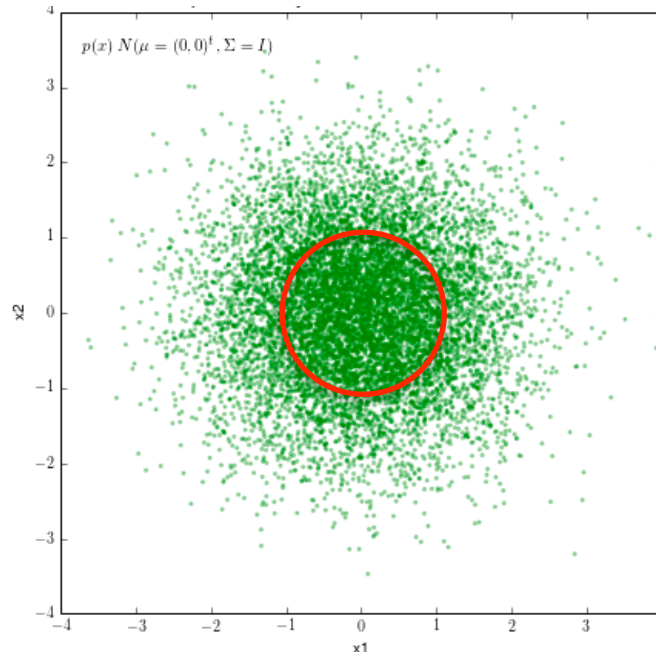
- However, as the dimension is increased something different happens
- When $\sigma^2 = 1$, integrating the probability density over a unit sphere centered at the origin yields nearly zero mass since the volume of a unit sphere is negligible.
- In fact, one needs to increase the radius of the sphere to \sqrt{d} before there is a significant nonzero volume and hence a nonzero probability mass.
- If one increases the radius beyond \sqrt{d} , the integral ceases to increase, even though the volume increases, since the probability density is dropping off at a much higher rate. The natural scale for the Gaussian is in units of $\sigma\sqrt{d}$.

Some Facts

$$E(x_1^2 + x_2^2 + \cdots + x_d^2) = d E(x_1^2) = d\sigma^2.$$

The probability mass of a unit-variance Gaussian as a function of the distance from its center is given by $r^{d-1}e^{-r^2/2}$ times some constant normalization factor where r is the distance from the center and d is the dimension of the space. The probability mass function has its maximum at

$$r = \sqrt{d - 1},$$



Concentration of Mass for Gaussian

Theorem 2.11 For a d -dimensional, unit variance, spherical Gaussian, for any positive real number $\beta < \sqrt{d}$, all but $3e^{-\frac{\beta^2}{8}}$ of the mass lies within the annulus $\sqrt{d}-\beta \leq r \leq \sqrt{d}+\beta$.

